Machine-Guided Discovery of a Real-World Rogue Wave Model

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DREVEN THEORYTHEORYTHEORYTHEORYTHEORYTHEORYTHEORYTHEORYTHEORYTHEORYTHEORYTHEORYTHEORYTHEORY Big data and large-scale machine learning have had a profound impact on science and engineering, particularly in fields focused on forecasting and prediction. Yet, it is still not clear how we can use the superior pattern matching abilities of machine learning models for scientific** *discovery***. This is because the goals of machine learning and science are generally not aligned. In addition to being accurate, scientific theories must also be causally consistent with the underlying physical process and allow for human analysis, reasoning, and manipulation to advance the field. In this paper, we present a case study on discovering a new symbolic model for oceanic rogue waves from data using causal analysis, deep learning, parsimony-guided model selection, and symbolic regression. We train an artificial neural network on causal features from an extensive dataset of observations from wave buoys, while selecting for predictive performance and causal invariance. We apply symbolic regression to distill this black-box model into a mathematical equation that retains the neural network's predictive capabilities, while allowing for interpretation in the context of existing wave theory. The resulting model reproduces known behavior, generates well-calibrated probabilities, and achieves better predictive scores on unseen data than current theory. This showcases how machine learning can facilitate inductive scientific discovery, and paves the way for more accurate rogue wave forecasting.** 1 $\overline{2}$ 3 4 5 6 7 8 9 10 11 12 13 14 15 16 17 18 19 20 21 22

ocean waves | rogue waves | machine learning | symbolic regression | causality

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gue waves are extreme ocean waves that have caused countless accidents, often with fatal consequences (1) \bullet countless accidents, often with fatal consequences (1). ³ They are defined as waves whose crest-to-trough height *H* ⁴ exceeds a threshold relative to the significant wave height ⁵ *Hs*. The significant wave height is defined as four times the standard deviation of the sea surface elevation. Here, we use ⁷ a rogue wave criterion with a threshold of 2.0:

8 *H/H_s* > 2.0 [1]

 A rogue wave is therefore by definition an unlikely sample from the tail of the wave height distribution, and can in principle occur by chance under any circumstance. This makes them difficult to analyze, and requires massive amounts of data. Therefore, research has mostly focused on theory and idealized experiments in wave tanks, often considering only 1-dimensional wave propagation [\(2\)](#page-11-1). However, the availability of large observation arrays [\(3\)](#page-11-2) makes them an ideal target for machine-learning based analysis [\(4,](#page-11-3) [5\)](#page-11-4).

 In this study, we present a neural network-based model that predicts rogue wave probabilities from the sea state, trained solely on observations from buoys (6) . The resulting model respects the causal structure of rogue wave generation; therefore, it can generalize to unseen physical regimes, is robust to distributional shift, and can be used to infer the relative importance of rogue wave generation mechanisms.

While a causally consistent neural network is useful for 25 prediction and qualitative insight into the physical dynamics, ²⁶ the ability for scientists to analyze, test, and manipulate a 27 model is crucial to recognize its limitations and integrate it 28 into the research canon. Despite advances in interpretable AI ²⁹ (7) , this is still a major challenge for most machine learning \sim 30 models. $\qquad \qquad \text{and} \qquad \qquad \text{and} \qquad \qquad \text{and} \qquad \$

To address this, we transform our neural network into a 32 concise equation using symbolic regression $(8, 9)$ $(8, 9)$ $(8, 9)$. The resulting $\frac{33}{2}$ model combines several known wave dynamics, outperforms ³⁴ current theory in predicting rogue wave occurrences, and can ³⁵ be interpreted within the context of wave theory. We see this 36 as an example of "data-mining inspired induction" (10) , an σ extension to the scientific method in which machine learning 38 guides the discovery of new scientific theories.

We achieve this through the following recipe (Fig. [1\)](#page-1-0): 40

- 1. A-priori analysis of causal pathways that leads to a set of ⁴¹ presumed causal parameters (Section [1\)](#page-0-0). ⁴²
- 2. Training an ensemble of regularized neural network pre- ⁴³ dictors, and parsimony-guided model selection based on ⁴⁴ causal invariance (Section [2\)](#page-2-0). ⁴⁵
- 3. Distillation of the neural network into a concise mathe- ⁴⁶ matical expression via symbolic regression (Section [3\)](#page-4-0). 47

Finally, we analyze both the neural network and symbolic ⁴⁸ model in the context of current wave theory (Section [4\)](#page-5-0). Both 49 models reproduce well-known behavior and point towards 50 new insights regarding the relative importance of different 51 mechanisms in the real ocean.

Significance Statement

Machine learning has had a transformative impact on predictive science and engineering. But due to their black-box nature, better machine learning models do not always lead to greater human understanding, the first goal of science. We show how this can be overcome by using machine learning to transform a vast database of wave observations into a new, human-readable equation for the occurrence probability of rogue waves — rare ocean waves that routinely damage ships and offshore structures. This equation can be analyzed and incorporated into the research canon. Our work demonstrates the potential of causal analysis, machine learning, and symbolic regression to drive scientific discovery in a real-world application.

All authors conceived the project. D.H. drafted, implemented, and executed the analysis. All authors interpreted the results. D.H. drafted the manuscript. All authors reviewed the manuscript

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Fig. 1. Overview of our study. Starting out with large amounts of tabular data from wave buoys, we use a causal analysis to identify the most important features for predicting rogue waves. We then train an ensemble of neural networks on subsets of these features, and select the best one based on its predictive performance and causal invariance. Finally, we use symbolic regression to distill the model into a concise mathematical equation. We analyze the neural network and symbolic expression in terms of their performance on unseen data, and compare them to existing theory. This closes the arc between data, machine learning, and theory.

⁵³ **1. A causal graph for rogue wave generation**

 To create a causal machine learning model it is crucial to expose it only to parameters with causal relevance. Otherwise, the model may prefer to encode spurious associations over true causal relationships, simply because they can be easier to learn. This requires us to identify the causal structure of rogue wave generation.

 There are several hypothesized causes of rogue waves (see [11,](#page-11-11) for an overview). Typically, research focuses on linear su- $_{62}$ perposition in finite-bandwidth seas (12) , wave breaking (13) , and wave-wave interactions in weakly nonlinear seas [\(14,](#page-11-14) [15\)](#page-11-15) or through the modulational instability (16) . Apart from these universal mechanisms, there are also countless possible interac- tions with localized features such as non-uniform topography [\(17\)](#page-11-17), wave-current interactions like in the Agulhas [\(18\)](#page-11-18) or the Antarctic Circumpolar Current [\(19\)](#page-11-19), or crossing sea states at high crossing angles affecting wave breaking (20) . We call this set of mechanisms the *physical effects* Φ.

 Since ocean waves are generated by a complex dynamical system, their true cause is a set of extrinsic *environmental conditions E* that are high-dimensional and not feasible to capture in full detail. However, most physical effects are mediated by one or several *sea state parameters* P, which are the characteristic aggregated parameters that appear in theoretical models of the respective wave dynamics, and that are included in operational wave forecasts. In this study, we would like obtain a model that relates relevant sea state 80 parameters P to wave observations O , which ideally also lets us infer the relative importance of physical effects Φ.

The go-to tool to analyze causal relationships is a causal 82 DAG (Directed Acyclic Graph; [21\)](#page-11-21). In a causal DAG, nodes 83 represent variables and edges $A \rightarrow B$ imply that *A* is a cause 84 of B (usually in the probabilistic sense in that the probability \qquad 85 distribution $P(B)$ depends on *A*).

We create a causal graph for rogue wave formation based 87 on the hypothesized causal mechanisms discussed above and set their corresponding theoretical models and parameters $(Fig. 2)$ $(Fig. 2)$. 89 Following this causal structure, we use the following set of sea 90 state parameters as candidates for representing the various 91 causal pathways (see Methods section for more information 92 on each parameter): 93

- Crest-trough correlation r , to account for the linear effect $\frac{1}{2}$ of wave groups on crest-to-trough rogue waves (22) . *r* 95 is the dominant causal factor behind linear rogue wave 96 formation (4) .
- Steepness *ε* governing weakly nonlinear effects, such as 98 second-order and third-order bound waves, and wave 99 $\frac{13}{100}$ [\(13,](#page-11-13) [23\)](#page-11-23).
- Relative high-frequency energy E_h (fraction of total energy contained in the spectral band 0*.*25 Hz to 1*.*5 Hz) as ¹⁰² a proxy for the strength of local winds (24) .
- Relative depth \overline{D} (based on peak wavelength), which is the central for nonlinear shallow-water effects (25, 26) and the central for nonlinear shallow-water effects $(25, 26)$ $(25, 26)$ $(25, 26)$ and wave breaking (13) .
- Dominant directional spread σ_{θ} , which has an influence 107

Fig. 2. The causes of rogue waves as a causal DAG (directed acyclic graph). Arrows $A \rightarrow B$ imply that A causes B .

¹⁰⁸ on third-order nonlinear waves [\(26\)](#page-11-26) and wave breaking $(20).$ $(20).$

• Spectral bandwidth ν_f (narrowness) and σ_f (peakedness), 111 appearing for example in the expression for the influence ¹¹² of third-order nonlinear waves (26).

¹¹³ We also include a number of derived parameters that commonly ¹¹⁴ appear in wave models and govern certain nonlinear (wave-¹¹⁵ wave) phenomena:

- ¹¹⁶ Benjamin-Feir index BFI, which controls third-order non- 117 linear free waves (26) and the modulational instability 118 (27) .
- ¹¹⁹ Ursell number Ur, which quantifies nonlinear effects in $_{120}$ shallow water (28) .

 \bullet Directionality index *R* (the ratio of directional spread and spectral bandwidth), which has an influence on third-order nonlinear free waves and is typically used in conjunction with the BFI (26) .

 These parameters cover most causal pathways towards rogue wave generation. Still, there are some at least partially unob- served causes, as we do not have access to data on local winds, topography, or currents. Additionally, our in-situ measure- ments are potentially biased estimates of the true sea state parameters, and there is no guarantee that any given training procedure will converge to the true causal model. This implies that we cannot rely on a model being causally consistent by design; instead, we perform a-posteriori verification on the learned models to find the perfect trade-off between causal consistency and predictive performance (see Section [2](#page-2-0)[C\)](#page-3-0).

¹³⁶ **2. An approximately causal neural network**

 A. Input data. We use the Free Ocean Wave Dataset (FOWD, [6\)](#page-11-5), which contains 1.4 billion wave measurements recorded by the 158 CDIP wave buoys [\(3\)](#page-11-2) along the Pacific and Atlantic coasts of the US, Hawaii, and overseas US territories. Water depths range between 10 m to 4000 m, and we require a sig- nificant wave height of at least 1 m. Each buoy records the sea surface elevation at a sampling frequency of 1*.*28 Hz, pro-ducing over 700 years of time series in total. FOWD extracts every zero-crossing wave from the surface elevation data and 145 computes a number of characteristic sea state parameters from ¹⁴⁶ the history of the wave within a sliding window.

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n certain no Due to the massive data volume of the full FOWD catalogue 148 $(\sim 1 \text{ TB})$, we use an aggregated version that maps each sea 149 state to the maximum wave height of the following 100 waves 150 (as in 4). This reduces the data volume by a factor of 100 and $_{151}$ inflates all rogue wave probabilities to a bigger value \hat{p} . We 152 correct for this via $p = 1 - (1 - \hat{p})^{1/100}$, assuming that rogue 153 waves occur independently from each other. This is a good 154 approximation in most conditions, but may underestimate ¹⁵⁵ seas with a strong group structure (see Section [5](#page-7-0)[B\)](#page-7-1).

The final dataset has $12.9M$ data points containing over 157 100*,*000 rogue waves exceeding 2 times the significant wave ¹⁵⁸ height. Our dataset is freely available for download (see Data 159 Availability section).

B. Neural network architecture. The probability to measure a 161 rogue wave based on the sea state can be modelled as a sum of 162 nonlinear functions, each of which only depends on a subset of 163 the sea state parameters representing a different causal path ¹⁶⁴ (act via different *physical effects* in Fig. [2\)](#page-2-1):

$$
logit P(y = 1 | \mathbf{x}) \sim \sum_{i} f_i(\mathbf{x}^{(S_i)}) + b
$$
 [2]

Here, ψ is a binary label indicating whether the current wave ψ is a rogue wave, $\mathbf{x}^{(S_i)}$ is the i-th subset of all causal sea state 168 parameters **x**, $logit(p) = log(p) - log(1-p)$ is the logit function, 169 f_i are arbitrary nonlinear functions to be learned, and *b* is a 170 constant bias term.

By including only a subset $\mathbf{x}^{(S_i)}$ of all parameters **x** as 172 input for f_i , we can restrict which parameters may interact $\frac{173}{2}$ non-additively with each other, which is an additional regu- ¹⁷⁴ larizing constraint that increases interpretability and prevents 175 interactions between inputs from different causal pathways. ¹⁷⁶ For example, to include the effects of linear superposition 177 and nonlinear corrections for free and bound waves (as in [29\)](#page-11-29), ¹⁷⁸ Eq. (2) can be written as: 179

$$
logit P(y = 1 | \mathbf{x}) \sim \underbrace{f_1(r)}_{linear} + \underbrace{f_2(BFI, R)}_{free \text{ waves}} + \underbrace{f_3(\varepsilon, \widetilde{D})}_{bound \text{ waves}} \quad [3] \quad 180
$$

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Fig. 3. Neural network architecture (multi-head FCN) used to predict rogue wave probabilities. Each input head receives a different subset of the full parameter set **x** to limit the amount of non-causal interactions between parameters.

 We use a neural network with fully-connected layers (FCN) to model the functions f_i , which are universal function ap- proximators [\(30\)](#page-11-30), and that can be trained efficiently for large ¹⁸⁴ amounts of data. The set of functions f_i can be represented as a single multi-head FCN with a linear output layer (Fig. 3). We use a small feed-forward architecture with 3 hidden lay-187 ers and ReLU activation functions (rectified linear units, 31). To the best of our knowledge, this is the first time that a multi-head neural network has been used to restrict the in- teractions between input parameters to be consistent with a causal model.

192 The neural network outputs a scalar $\tilde{p} \in (-\infty, \infty)$, the log-¹⁹³ odds of a rogue wave occurrence for the given sea state. For 194 training, we use the Adam optimizer (32) and backpropagation ¹⁹⁵ to minimize a cross-entropy loss for binary classification with 196 an added ℓ_2 regularization term for kernel parameters:

$$
L(p, y, \theta) = y \cdot \log(p) + (1 - y) \cdot \log(1 - p) + \lambda \|\theta\|_2 \quad [4]
$$

with predicted probability $p = logit^{-1}(\tilde{p})$, observed labels 199 $y \in \{0,1\}$ (rogue wave or not), and neural network kernel ²⁰⁰ parameters *θ*.

 To estimate uncertainties in the neural network parameters and resulting predictions, we use Gaussian stochastic weight averaging (SWAG, [33\)](#page-11-33). For this, we train the network for 50 epochs, then start recording the optimizer trajectory after each epoch for another 50 epochs. The observed covariance structure of the sampled parameters is used to construct a multivariate Gaussian approximation of the loss surface that we can sample from. This results in slightly better predictions, and gives us a way to quantify how confident the neural 209 network is in its predictions. 210

C. Causal consistency and predictive accuracy. Although we 211 include only input parameters that we assume to have a di- ²¹² rect causal connection with rogue wave generation, there is 213 no guarantee that the neural network will infer the correct ²¹⁴ causal model. In fact, the presence of measurement bias and ²¹⁵ unobserved causal paths makes it unlikely that the model will ²¹⁶ converge to the true causal structure. To search for an approx- ²¹⁷ imately causally consistent model we will have to quantify its ²¹⁸ causal performance.

We achieve this through the concept of invariant causal 220 prediction (ICP; 34 , 35). The key insight behind ICP is that 221 the parameters of the true causal model will be invariant under 222 distributional shift, that is, an intervention on an upstream 223 "environment" node in the causal graph that controls which ²²⁴ distribution the data is drawn from. Re-training the model 225 on data with different spurious correlations *between* features 226 should still lead to the same dependency of the target *on* the 227 features (see also 36).

We split the dataset randomly into separate training and 229 validation sets, in chunks of 1M waves. We train the model on 230 the full training dataset and perform ICP on the validation 231 dataset, which we partition into subsets representing different conditions in space, time, depth, spectral properties, and 233 degrees of non-linearity (Table [1\)](#page-4-1). This changes the domi- ²³⁴ nant characteristics of the waves in each subset (representing 235 e.g. storm and swell conditions), inducing distributional shift. ²³⁶ Then, we re-train the model separately on each subset and 237 compute the root-mean-square difference between predictions 238

Table 1. The subsets of the validation dataset used to evaluate model performance and invariance.

Subset name	Condition	$#$ waves
southern-california	Longitude $\in (-123.5, -117)^{\circ}$, latitude $\in (32, 38)^{\circ}$	265M
deep-stations	Water depth $>1000 \text{ m}$	28M
shallow-stations	Water depth $< 100 \,\mathrm{m}$	154M
summer	Day of year $\in (160, 220)$	51M
winter	Day of year $\in (0, 60)$	91M
$\text{Hs} > 3\text{m}$	$H_s > 3m$	58M
high-frequency	Relative swell energy < 0.15	43M
low-frequency	Relative swell energy > 0.7	46M
long-period	Mean zero-crossing period $> 9s$	100M
short-period	Mean zero-crossing period $< 6s$	42M
cnoidal	Ursell number > 8	40M
weakly-nonlinear	Steepness > 0.04	83M
low-spread	Directional spread $< 20^{\circ}$	25M
high-spread	Directional spread $>40^{\circ}$	25M
full	(all validation data)	472M

239 of the re-trained model P_k and the full model P_{tot} on the k -th 240 data subset $\mathbf{x}_{(k)}$:

$$
\mathcal{E}_k^2 = \frac{1}{n_k} \sum_{i}^{n_k} \left(\text{logit } P_k\left(\mathbf{x}_i^{(k)}\right) - \text{logit } P_{\text{tot}}\left(\mathbf{x}_i^{(k)}\right) \right)^2 \quad [5]
$$

²⁴² As the total consistency error we use the root-mean-square of ²⁴³ Eq. [\(5\)](#page-4-2) across all environments:

$$
\mathcal{E} = \sqrt{\frac{1}{n_E} \sum_{k}^{n_E} \mathcal{E}_k^2}
$$
 [6]

 Under a noise-free, infinite dataset and an unbiased training process that always identifies the true causal model we would ²⁴⁷ find $\mathcal{E} = 0$, i.e., re-training the model on the unseen data subset would not contribute any new information and leave the model perfectly invariant. Since all of these assumptions are violated here, we merely search for an approximately causal 251 model that minimizes \mathcal{E} .

252 However, we cannot use $\mathcal E$ as the only criterion when se- lecting a model. The invariance error can only account for change in the prediction (variance), but not for its overall closeness to the true solution (bias). Therefore, we select a model that is Pareto-optimal with respect to the invariance ϵ_{257} error $\mathcal E$ and a predictive score $\mathcal L$. This will not establish ab- solute causal consistency, but will allow us to select a model that is near-optimal given the constraints.

 ϵ_{260} For $\mathcal L$ we use the log of the likelihood ratio between the ²⁶¹ predictions of our neural network and a baseline model that predicts the empirical base rate $\overline{y}_k = \frac{1}{n} \sum_{i=1}^{n} y_{k,i}$, averaged over ²⁶³ all environments *k*:

264
$$
\mathcal{L}(p,\overline{y}) = \frac{1}{n_E} \sum_{k}^{n_E} (I(p_k) - I(\overline{y}_k))
$$
 [7]

$$
I(x) = x \cdot \log(x) + (1 - x) \cdot \log(1 - x) \tag{8}
$$

²⁶⁶ To evaluate model calibration (the tendency to produce over-²⁶⁷ or under-confident probabilities), we compute a calibration curve by binning the predicted rogue wave probabilities. We ²⁶⁸ then compare each bin to the observed rogue wave frequency, 269 and compute the weighted root-mean-square residual between 270 measured (\overline{y}_i) and predicted (p_i) log-odds: 271

$$
C = \sqrt{\sum_{i=1}^{n_b} w_i \left(\text{logit}(p_i) - \text{logit}(\overline{y}_i) \right)^2}
$$
 [9] ₂₇₂

To account for uncertainty in the observations (e.g. close to ²⁷³ the extremes), the weights w_k are based on the 33% credible 274 \int interval of \overline{y}_i ∼ Beta (n_i^+, n_i^-) with n_i^+ rogue and n_i^- non-rogue 275 measurements. This is similar to the expected calibration error ²⁷⁶ (37) , but models data uncertainty directly. We use a uniform 277 bin size (in logit space) of 0.1 .

D. Model selection. We train a total of 24 candidate models on 279 different subsets of the relevant causal parameters (as identified 280 in Section [1\)](#page-0-0) and varying number of input heads (between 1 281 and 3). We evaluate their performance in terms of calibration, ²⁸² predictive performance, and causal consistency (Table [3\)](#page-10-0). 283

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full model P_{tot} on the k-th
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the other hand (Fig. 4). [T](#page-5-1)his is
models are indeed less biased
connections. We perform model
 $-\logit P_{\text{tot}}(\mathbf{x}_i^{(k)})\big)^2$ [5 We observe a clear anti-correlation between model complex- ²⁸⁴ ity and predictive score on one hand and causal consistency on ²⁸⁵ the other hand (Fig. 4). This is evidence that more complex 286 models are indeed less biased but exploit more non-causal ²⁸⁷ connections. We perform model selection based on *parsimony*: ²⁸⁸ A good model is one where a small increase in either predictive ²⁸⁹ performance or causal consistency implies a large decrease in ²⁹⁰ the other, i.e., where the Pareto front is convex. This is similar ²⁹¹ to the metric used by PySR (9) to select the best symbolic 292 regression model (Section 3).

Based on this, we choose model 18 with parameter groups ²⁹⁴ $S_1 = \{r\}, S_2 = \{\varepsilon, \sigma_\theta, \sigma_f, D\}$ (i.e., a model with two input 295 heads) as the reference model for further analysis. The chosen heads) as the reference model for further analysis. The chosen model produces well-calibrated probabilities (Fig. 5), and is 297 among the 5 best models in terms of predictive performance ²⁹⁸ on the test dataset (not used during training or selection), ²⁹⁹ despite using only 5 features with at most 4-way interactions. 300

The relatively low number of input features allows us to 301 analyze the model in detail using explainable AI methods ³⁰² $(Section 4A)$. 303

3. Learning an empirical equation for rogue wave risk ³⁰⁴

To make our model fully interpretable, we transform the ³⁰⁵ learned neural network into an equation via symbolic regres- ³⁰⁶ sion. Common approaches to symbolic regression include 307 Eureqa (39) , AI Feynman (40) , SINDy (41) , and QLattice 308 (42) . Here, we use PySR $(8, 9)$ $(8, 9)$ $(8, 9)$, a symbolic regression pack- 309 age based on genetic programming (43) . Genetic algorithms $\frac{310}{2}$ build a large ensemble of candidate models and select the ³¹¹ best ones, before mutating and recombining them into the ³¹² next generation. In the case of symbolic regression, mathe- ³¹³ matical expressions are represented as a tree of constants and 314 elementary symbols. In principle, this allows $P_{\rm V}$ SR to discover $\frac{315}{2}$ expressions of unbounded complexity. 316

PySR's central metric to quantify the goodness of an equa- 317 tion is again based on *parsimony*, in the form of the derivative 318 of predictive performance with respect to the model complex- ³¹⁹ $ity - if the true model has been discovered, any additional$ complexity can at best lead to minor performance gains (by ³²¹ $overfitting to noise in the data).$

Fig. 4. There is a clear trade-off between causal invariance (\mathcal{E}) and predictive performance (\mathcal{E}) of our neural network predictors. We choose the model that lies in the most convex part of the Pareto frontier. Scores are evaluated on validation data. Test performance is based on prediction scores on held-out test data (from unseen stations).

Fig. 5. Our model outputs well-calibrated probabilities, even for unseen stations. Shown is the binned predicted probability *p* vs. the observed rogue wave frequency \overline{y} on the test data. Error bars for p indicate 3 standard deviations estimated via SWAG sampling. Error bars for \overline{y} indicate 95 % credible interval assuming $\overline{y}_i \sim$ $\mathrm{Beta}(n^+_i,n^-_i)$. Bins with less than 10 observed rogue waves are excluded. Dashed line indicates perfect calibration. Solid line indicates probability as predicted by linear theory in the narrow-bandwidth limit (Rayleigh distribution; 38).

 In our case, we seek to find an expression *f* from the space of possible expression graphs T*^O* with allowed operators *O* that approximates the rogue wave log-probability as predicted 326 by the neural network $\mathcal N$ over the dataset x :

$$
\text{Find } f \in \mathcal{T}_O \text{ that minimizes } \sum_i \frac{1}{\text{Var}(y_i)} \left[f(x_i)^2 - \sigma(\mathbb{E}[y_i])^2 \right]
$$

328 where $\sigma(x) = -\log(1 + \exp(-x))$, and y_i is the set of SWAG 329 samples from $\mathcal{N}(x_i)$. A sensible set of operators O is key to ³³⁰ ensure interpretability of the resulting expression; we choose the symbols $O = \{+, -, \times, \div, \log, \{-1, \sqrt{2}\} \}$ to facilitate ex-³³² pressions that are similar to current theoretical models of 333 the form $P \sim A \exp(B)$. We normalize all input features to ³³⁴ approximately unit scale by converting directional spread to ³³⁵ radians.

³³⁶ PySR assembles a league of candidate expression and ³³⁷ presents the Pareto-optimal solutions of increasing complexity ³³⁸ to the user. We select the best solution by hand, picking the expression with the best parsimony score that contains all in- ³³⁹ put features and at least two terms containing the steepness ε 340 (to account for the various causal pathways in which steepness $_{341}$ affects rogue waves). The final equation is shown in Fig. 8 , $\frac{342}{2}$ and discussed in Section [4B.](#page-7-3) 343

4. Results 344

A. Neural network. We analyze the behavior of our neural 345 network predictor, which reveals important insights about the ³⁴⁶ physical dynamics of rogue waves and their prediction. $\frac{347}{2}$

A.1. Rogue wave models should account for crest-trough correlation, ³⁴⁸ steepness, relative depth, and directionality. Only this parameter 349 combination achieves good causal consistency and predictive 350 scores at the same time, and experiments that exclude any of 351 these parameters perform unconditionally worse in either met- ³⁵² ric. Especially the exclusion of crest-trough correlation leads 353 to catastrophic results, even when including other bandwidth 354 measures like σ_{θ} in its place (Table [3\)](#page-10-0).

This suggests that the above set of parameters represents 356 the dominant rogue wave generation processes in the form of 357 linear superposition in finite-bandwidth seas with a directional 358 contribution and weakly nonlinear corrections.

The crest-trough correlation r is still lacking mainstream $\frac{360}{250}$ adoption as a rogue wave indicator (for example, it is not part 361 of ECMWF's operational forecast; [29\)](#page-11-29), despite being a key ³⁶² parameter for crest-to-trough rogue waves $(4, 22, 44)$ $(4, 22, 44)$ $(4, 22, 44)$ $(4, 22, 44)$ $(4, 22, 44)$. The 363 other parameters are consistent with other empirical studies ³⁶⁴ such as Fedele (45) , which considers the same parameters in $\frac{365}{200}$ conjunction with rogue crests during storms. They are also 366 similar to the ingredients to ECMWF's rogue wave forecast $\frac{367}{267}$ (29) , which is based on the effects of second and third-order $\frac{1}{368}$ bound and free waves and uses steepness, relative depth, direc- ³⁶⁹ tional spread, and spectral bandwidth. However, in our model ³⁷⁰ these parameters are combined differently; a model enforcing 371 the same interactions (steepness and relative depth for bound 372 wave contribution, BFI and directionality index for free wave 373 contribution) performs poorly. 374

Numerous previous studies have found the BFI to be a 375 poor predictor of rogue wave risk in realistic sea states [\(4,](#page-11-3) ³⁷⁶ [14,](#page-11-14) [15,](#page-11-15) [45](#page-11-45)[–48\)](#page-11-46) due to its strong underlying assumptions such $\frac{377}{27}$ as unidirectionality. This study extends this to the fully $\frac{378}{276}$ nonparametric and nonlinear case.

Fig. 6. ALE (accumulated local effects) plot matrix for experiment 18. Shown is the change in rogue wave risk (in logits) from the average as each parameter is varied. The total effect is the sum of all 1D, 2D, and higher-order contributions (not shown).

 We study how our model uses different parameters by visu- alizing their impact on the prediction of the respective head of the neural network. For this, we make use of the accumulated local effects decomposition (ALE, 49), which measures the influence of infinitesimal changes in each parameter on the 385 prediction outcome (see also 7). From the ALE plot (Fig. 6), we find that crest-trough correlation has by far the biggest influence of all parameters and explains about 1 order of mag- nitude in rogue wave risk variation, which is consistent with earlier model-free approaches [\(4\)](#page-11-3). To first order, higher crest- trough correlation, lower directional spread, larger relative depth (deep water), and higher steepness lead to larger rogue wave risk, but parameter interactions can lead to more com- plicated, non-monotonic relationships (for example in very shallow water, see Section [4](#page-5-0)[A.3\)](#page-6-1).

 A.2. The Rayleigh distribution is an upper bound for real-world rogue wave risk. Despite the clear enhancement by weakly nonlinear corrections, the Rayleigh wave height distribution remains an upper bound for real-world (crest-to-trough) rogue waves. The Rayleigh distribution is the theoretical wave height distribution 400 for linear narrow-band waves [\(38\)](#page-11-43), i.e., the limit $r \to 1$, $\varepsilon \to 0$, $\sigma_f \to 0$, $\overline{D} \to \infty$, and $\sigma_\theta \to 0$, and reads:

$$
P(H/H_s > k) = \exp(-2k^2)
$$
 [10]

Fig. 7. Our model predicts a positive association between steepness and rogue waves in deep water, and a negative association in shallow water. Shown is the 1-dimensional ALE (accumulated local effects) plot in both cases. Here, deep water are sea states with $\widetilde{D} > 3$ and shallow water with $\widetilde{D} < 0.1$.

Only in the most extreme conditions does our model predict a 403 similarly high probability, for example for $\sigma_{\theta} = 13^{\circ}$, $\varepsilon = 0.008$, 404 $\sigma_f = 0.14$, $r = 0.88$, and $D = 0.6$, which gives the same 405 probability as the Rayleigh distribution, $p = 3.3 \times 10^{-4}$. . ⁴⁰⁶

In the opposite extreme, rogue wave probabilities can fall $_{407}$ to as little as 10^{-5} for low values of *r* and high values of σ_{θ} 408 (such as in a sea with a strong high-frequency component and ⁴⁰⁹ high directional spread). This suggests that bandwidth effects $\frac{410}{2}$ can create sea states that efficiently suppress extremes. ⁴¹¹

$$
P(H > 2H_S \mid r, \varepsilon, \sigma_\theta, \sigma_f, \widetilde{D}) = \exp \left[-12. + 3.8r - \frac{\log \sigma_\theta}{2} + 66. \varepsilon^2 - \sqrt{\varepsilon} - \underbrace{\frac{0.23\varepsilon}{\widetilde{D} \cdot \sigma_f}}_{\text{II}} \right]
$$

Fig. 8. Our empirical equation for roque wave risk, as identified through the distillation of our neural network predictor via symbolic regression. This equation outperforms existing wave theory on unseen stations from our dataset, while being fully interpretable. Numbered terms are discussed in Section [4](#page-5-0)[B.](#page-7-3) All floating point coefficients are rounded to two significant digits.

 A.3. There is a clear separation between deep water and shallow water regimes. All models with high causal invariance scores include an interaction between steepness and relative water depth. Looking at this more closely, we find that a stratification on deep and shallow water sea states reveals 2 distinct regimes 417 (Fig. [7\)](#page-6-2).

 In deep water, rogue wave risk is strongly positively as- sociated with steepness, as expected from the contribution of second and third-order nonlinear bound waves (26) . The 421 opposite is true in shallow water $(D < 0.1)$, where we find a clear *negative* association with steepness. This is likely due clear *negative* association with steepness. This is likely due to depth-induced wave breaking [\(23\)](#page-11-23). In very shallow waters, more sea states have a steepness close to the breaking thresh- old, which removes taller waves that tend to have a higher steepness than average.

 B. Symbolic expression. The final expression for the rogue wave probability, as discovered via symbolic regression, is given in Fig. [8.](#page-7-2) It consists of an exponential containing five additive ⁴³⁰ terms:

 $_{431}$ (I) $-12 + 3.8r$. The term with the largest coefficients is the 432 one containing r , as expected. Comparison with the ex-433 ponential term in the Tayfun distribution P_t , Eq. (28), re-⁴³⁴ veals that this is approximately a linear expansion around 435 $r \approx 1$:

$$
\log P_t(H/H_s > h) \sim -\frac{4h^2}{1+r} \tag{11}
$$

$$
= -12 + 4r + \mathcal{O}(r^2) \Big|_{r \approx 1}^{n-2} \qquad [12]
$$

⁴⁴⁰ from data.

⁴³⁸ This is an important sanity check for the model, since it ⁴³⁹ shows that it is able to re-discover existing theory purely

 $= -12 + 4r + \mathcal{O}(r^2) \Big|_{r \approx 1}^{h=2}$

- 441 (II) $-\log \sigma_{\theta}/2$. This encodes the observed enhancement for ⁴⁴² narrow sea states and has no direct relation to existing ⁴⁴³ quantitative theory. Its functional form is somewhat ⁴⁴⁴ problematic, since it causes the model to diverge for $\sigma_{\theta} \rightarrow 0$ (unidirectional seas). However, the model has 446 only seen real-world seas with $\sigma_{\theta} \gtrsim 0.2$, so we may replace ⁴⁴⁷ this term with one that yields similar predictions for the relevant range of σ_{θ} , and does not diverge for $\sigma_{\theta} \to 0$.
- ⁴⁴⁹ One possible candidate is

$$
\frac{1 - \sigma_{\theta}}{1 + \sigma_{\theta}}, \tag{13}
$$

- ⁴⁵¹ which has a relative RMS error of about 5 % over the ⁴⁵² range $\sigma_{\theta} \in (20, 90)^{\circ}$ compared to the original term.
- 453 (III) $66\varepsilon^2$. Encodes the influence of weakly nonlinear effects 454 for large values of $\varepsilon \geq 0.1$.
- (IV) $-\sqrt{\varepsilon}$. This term encodes the observed negative association 455 between steepness and rogue waves for low values of ε 456 that could be due to wave breaking, or may be an artifact 457 of our sensor. 458
- (V) $0.23\varepsilon/(\overline{D}\cdot\sigma_f)$. Since $\overline{D}\sim k_pD$ and $\varepsilon\sim k_pH_s$, this term 459 is proportional to the relative wave height $\eta = H_s/D$ 460 is proportional to the relative wave height $\eta = H_s/D$ and $1/\sigma_f$. *η* is the most important parameter in the 461 theory of shallow-water waves, and appears for example 462 in the Korteweg-de Vries equation (25) . Accordingly, 463 this term dominates the dynamics in very shallow water. ⁴⁶⁴ Dependencies on $1/\sigma_f$ occur in current theory [\(26\)](#page-11-26), but \sim 465 are usually paired with σ_{θ} to form the directionality index $\frac{466}{4}$ *R*. This suggest that term V may be incomplete, and $\frac{467}{467}$ missing physical dynamics that are not prevalent in the ⁴⁶⁸ $data.$ 469

Overall, the equation is able to reproduce the same qualitative behavior as observed from the neural network, with the same well-calibrated outputs ($\mathcal{C} = 0.14$) and predictive performance 472 $(Section 5A)$ on the test data.

5. Discussion ⁴⁷⁴

2.5). In very shahow waters,

lose to the breaking thresh-

that tend to have a higher

missing physical dynamic

data.

al expression for the rogue

overall, the equation is able to

symbolic regression, is given

that c **A. Validation against theory.** We test our models (neural net- ⁴⁷⁵ work and symbolic equation) against existing wave theory ⁴⁷⁶ based on their mean predictive score $\mathcal L$ across the environments 477 from Table 1 on the held-out test data (unseen stations). As ⁴⁷⁸ theoretical baselines we use the models from Longuet-Higgins 479 (Rayleigh, 38), Tayfun (22), Mori & Janssen [\(50\)](#page-11-48), and a hybrid $\frac{480}{90}$ combining Tayfun and Mori & Janssen (see Methods section). ⁴⁸¹

The results are shown in Fig. [9.](#page-8-0) Since the Rayleigh and 482 Mori & Janssen models do not account for crest-trough correlation, their predictions vastly overestimate the occurrence ⁴⁸⁴ rate of observed rogue waves. The Tayfun and hybrid models 485 perform better, but are still outperformed by our models ex- ⁴⁸⁶ cept in cnoidal seas. Our models are better predictors than ⁴⁸⁷ the baseline (predicting the empirical per-environment rogue 488 wave frequency) in all environments.

The neural network performs better than the symbolic equa- ⁴⁹⁰ tion in all environments, albeit only by a small margin. This ⁴⁹¹ shows that the symbolic equation is able to capture the main 492 features of the full model, despite its compact representation. ⁴⁹³

B. Limitations. Using only wave buoy observations for our 494 analysis, we acknowledge the following limitations: 495

• We did not have sufficient data on local winds, currents, ⁴⁹⁶ or topography, which implies that some relevant causal ⁴⁹⁷ pathways are unobserved (see Fig. [2\)](#page-2-1). While we expect ⁴⁹⁸ these effects to play a minor role in bulk analysis, they 499 could dramatically affect local rogue wave probabilities in 500 specific conditions, for example over sloping topography 501 (17) or in strong currents (51) .

Fig. 9. Comparison between our models and existing theory on held-out test data. Our models perform similar to each other and outperform existing theory on this dataset in all but one data subset (cnoidal seas). *x*-scale is linear in $(-10^{-3}, 10^{-3})$, and logarithmic otherwise.

 • We only have one-dimensional (time series) data and cannot capture imported parameters, such as solitons generated elsewhere that travel into the observation area. While we expect this to play a minor role, it could under-estimate the importance of nonlinear free waves.

 • Systematic sensor bias is common in buoys and can lead to spurious causal relationships. This may obscure the true causal structure and hurt model generalization to other sensors. However, this adaptation to sensor characteristics may be desirable in forecasting scenarios, where it allows the model to synthesize several noisy quantities into more robust ones.

 • By aggregating individual waves into 100-wave chunks, we underestimate the per-wave rogue wave probability in sea states in which rogue waves do not occur independently of each other, such as seas with a strong group structure.

 These limitations could potentially reduce our model's ability to detect relevant causal pathways and underestimate the true rogue wave risk. Our analysis is agnostic to the data source and can be repeated on different sources to validate our findings.

⁵²⁴ **6. Next steps**

 A. An improved rogue wave forecast. Our empirical model can be compared directly to existing rogue wave risk indicators by evaluating them on forecast sea state parameters. ECMWF's operational rogue wave forecast [\(29\)](#page-11-29) focuses on envelope wave heights which does not account for crest-trough correlation, and is conceptually similar to the Mori & Janssen model in Section [5](#page-7-0)[A.](#page-7-4) Therefore, we are confident that substantial improvements are within reach in terms of predicting crest-to-trough rogue waves, even without using a black-box model.

B. Predicting super-rogue waves. Observed wave height distri-
534 butions often show a flattening of the wave height distribution 535 towards the extreme tail $(11, 14, 52)$ $(11, 14, 52)$ $(11, 14, 52)$ $(11, 14, 52)$ $(11, 14, 52)$. Therefore, we expect $\overline{}$ 536 rogue wave probabilities to be more pronounced for even more 537 extreme waves (for example with $H/H_s > 2.4$, as recently 538 α observed in 53).

PRECIST: The set of the observation area. In the observation area The lack of sufficient direct observations in these regimes 540 calls for a different strategy. One approach could be to trans- ⁵⁴¹ form this classification problem (rogue wave or not) into a ⁵⁴² regression, where the predicted variables are the free parame- ⁵⁴³ ters of a candidate wave height probability distribution (such ⁵⁴⁴ as shape and scale parameters of a Weibull distribution). Then, ⁵⁴⁵ a similar analysis as in this study could be conducted for these ⁵⁴⁶ parameters, which may reveal the main mechanisms influenc- ⁵⁴⁷ ing the risk for truly exceptional waves, and whether this ⁵⁴⁸ flattening can be confirmed in our dataset. 549

> **C. Commoditization of data-mining based induction.** There is 550 a pronounced lack of established methods for machine learning 551 aimed at scientific discovery. We have shown that incorpo- ⁵⁵² rating and enforcing causal structure can overcome many of 553 the shortcomings of standard machine learning approaches, ⁵⁵⁴ like poorly calibrated predictions, non-interpretability, and 555 incompatibility with existing theory. However, the methods 556 we leveraged are still in their infancy and rely on further com- ⁵⁵⁷ munity efforts to be end-to-end automated and adopted at 558 scale. Particularly, parsimony-based model selection (as in 559 Section [2](#page-2-0)[D](#page-4-3) and Section [3\)](#page-4-0) is still a manual process that requires a firm understanding of model intrinsics and the domain 561 at hand. Nonetheless, we believe that the potential benefits of 562 causal and parsimony-guided machine learning for real-world 563 problems are too great to ignore, and we hope that this study ⁵⁶⁴ will inspire further research in this direction.

⁵⁶⁶ **Materials and Methods**

 Sea state parameters. Here, we give the definition of the sea state parameters used in this study. For a more thorough description of how parameters are computed from buoy displacement time series see Häfner et al. (6) .

 All parameters can be derived from the non-directional wave 572 spectrum $S(f)$, with the exception of directional spread $σ_θ$, which is estimated from the horizontal motion of the buoy and taken from the raw CDIP data.

⁵⁷⁵ Most parameters are computed from moments of the wave spec- 576 trum, where the *n*-th moment m_n is defined as

$$
m_n = \int_0^\infty f^n S(f) \, df \qquad [14]
$$

⁵⁷⁸ The expressions for the relevant sea state parameters are:

 $$

⁵⁷⁹ • Significant wave height:

$$
H_s = 4\sqrt{m_0} \tag{15}
$$

⁵⁸¹ • Spectral bandwidth (narrowness):

$$
\nu_f = \sqrt{m_2 m_0 / m_1^2 - 1} \tag{16}
$$

⁵⁸³ • Spectral bandwidth (peakedness):

$$
\sigma_f = \frac{m_0^2}{2\sqrt{\pi}} \left(\int_0^\infty f \cdot \mathcal{S}(f) \, \mathrm{d}f \right)^{-1} \tag{17}
$$

⁵⁸⁵ • Peak wavenumber *kp*, computed via the peak period (as in 586 54 :

$$
\overline{T}_p = \frac{\int S(f)^4 \, df}{\int f \cdot S(f)^4 \, df} \tag{18}
$$

⁵⁸⁸ This leads to the peak wavenumber through the dispersion ⁵⁸⁹ relation for linear waves in intermediate water of depth *D*:

590
$$
f(k)^2 = \frac{gk}{(2\pi)^2} \tanh(kD)
$$
 [19]

$$
591 \qquad \qquad \text{An approximate inverse is given in Fenton (55)}.
$$

592 • Relative depth, based on the wave length λ :

$$
\widetilde{D} = \frac{D}{\lambda} = \frac{1}{2\pi} k_p D \tag{20}
$$

⁵⁹⁴ • Peak steepness: 595 $\varepsilon = H_s k_p$ [21]

⁵⁹⁶ • Benjamin-Feir index:

$$
BFI = \frac{\varepsilon \nu}{\sigma_f} \sqrt{\max\{\beta/\alpha, 0\}} \tag{22}
$$

598 where *ν*, *α*, *β* are coefficients depending only on \widetilde{D} (full expression given in 56). sion given in 56).

⁶⁰⁰ • Directionality index:

$$
R = \frac{\sigma_{\theta}^2}{2\nu_f^2} \tag{23}
$$

⁶⁰² • Crest-trough correlation:

$$
r = \frac{1}{m_0} \sqrt{\rho^2 + \lambda^2} \tag{24}
$$

$$
\rho = \int_0^\infty \mathcal{S}(\omega) \cos\left(\omega \frac{\overline{T}}{2}\right) d\omega \qquad [25]
$$

$$
\lambda = \int_0^\infty \mathcal{S}(\omega) \sin\left(\omega \frac{\overline{T}}{2}\right) d\omega \qquad [26]
$$

606 where ω is the angular frequency and $\overline{T} = m_0/m_1$ the spectral ⁶⁰⁷ mean period [\(12\)](#page-11-12).

Model implementation and hyperparameters. All performance critical 608 model code is implemented in JAX (57) , using neural network 609 modules from flax (58) and optimizers from optax (59) . We run 610 each experiment on a single Tesla P100 GPU in about 40 minutes, 611 including SWAG sampling and re-training on every validation subset. ⁶¹² The whole training process can also be executed on CPU in about 613 2 hours. The hyperparameters for all experiments are shown in 614 $Table 2.$ $Table 2.$

Table 2. Hyperparameters used in experiments.

nh: number of input heads.

Full list of experiments. See Table [3.](#page-10-0) 616

s):
 [F](#page-11-43) $\cdot S(f) df$ $\left(\frac{17}{4} + \frac{1}{5} + \frac{1}{15}\right)$ **Fig. 117 IV**
 Example 12 Example 12 Example 12 Example 135 Full list of experiments. See Table
 Example 13 Example 14 Example 14 Example 14 Example 14 Reference wave height distributions. We use the following theoretical 617 wave height exceedance distributions for comparison (with rogue 618 wave threshold κ , here $\kappa = 2$): 619

• Rayleigh (38) : 620

$$
P_{\rm R}(\kappa) = \exp\left(-2\kappa^2\right) \tag{27}
$$

 $\arctan(12, 22)$: 622

$$
P_{\rm T}(\kappa) = \exp\left(\frac{-4}{1+r}\kappa^2\right) \tag{28}
$$

Mori & Janssen $(50, 60)$: 624

$$
P_{\rm MJ}(\kappa) = \left(1 + \frac{2\pi}{3\sqrt{3}} \frac{\rm BFI^2}{1 + 7.1R} \kappa^2 (\kappa^2 - 1)\right) \exp\left(-2\kappa^2\right) [29] \quad \text{625}
$$

• Hybrid: 626

$$
P_{\rm H}(\kappa) = \left(1 + \frac{2\pi}{3\sqrt{3}} \frac{\rm BFI^2}{1 + 7.1R} \kappa^2 (\kappa^2 - 1)\right) \exp\left(\frac{-4}{1 + r} \kappa^2\right) [30] \quad \text{e27}
$$

Data Availability. The preprocessed and aggregated version of the 628 FOWD CDIP data used in this study is available for download 629 [a](https://erda.ku.dk/archives/ee6b452c1907fbd48271b071c3cee10e/published-archive.html)t [https://erda.ku.dk/archives/ee6b452c1907fbd48271b071c3cee10e/](https://erda.ku.dk/archives/ee6b452c1907fbd48271b071c3cee10e/published-archive.html) 630 [published-archive.html](https://erda.ku.dk/archives/ee6b452c1907fbd48271b071c3cee10e/published-archive.html). All model code is openly available at 631 <https://github.com/dionhaefner/rogue-wave-discovery>. 632

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Table 3. Full list of experiments. L**: Prediction score (higher is better).** E**: Invariance error (lower is better).** C**: Calibration error (lower is better). Color coding ranges between** (**median** − **IQR***,* **median** + **IQR**) **with inter-quartile range IQR.**

- 648 1. E Didenkulova, Catalogue of rogue waves occurred in the World Ocean from 2011 to 2018 649 reported by mass media sources. *Ocean. & Coast. Manag*. p. 105076 (2019).
- 650 2. JM Dudley, G Genty, A Mussot, A Chabchoub, F Dias, Rogue waves and analogies in optics 651 and oceanography. *Nat. Rev. Phys*. **1**, 675–689 (2019) Number: 11 Publisher: Nature 652 Publishing Group.
- 653 3. J Behrens, J Thomas, E Terrill, R Jensen, CDIP: Maintaining a Robust and Reliable Ocean 654 Observing Buoy Network in *2019 IEEE/OES Twelfth Current, Waves and Turbulence Measure-*655 *ment (CWTM)*. pp. 1–5 (2019).
- 656 4. D Häfner, J Gemmrich, M Jochum, Real-world rogue wave probabilities. *Sci. Reports* **11**, 657 10084 (2021) Number: 1 Publisher: Nature Publishing Group.
- 658 5. A Cattrell, M Srokosz, B Moat, R Marsh, Can rogue waves be predicted using characteristic 659 wave parameters? *J. Geophys. Res. Ocean*. **123**, 5624–5636 (2018).
- 660 6. D Häfner, J Gemmrich, M Jochum, FOWD: A Free Ocean Wave Dataset for Data Mining 661 and Machine Learning. *J. Atmospheric Ocean. Technol*. **-1** (2021) Publisher: American 662 Meteorological Society Section: Journal of Atmospheric and Oceanic Technology.
- 663 7. C Molnar, *Interpretable Machine Learning*. (2020).
- 664 8. M Cranmer, et al., Discovering symbolic models from deep learning with inductive biases. 665 *NeurIPS 2020* (2020).
- 666 9. M Cranmer, Interpretable machine learning for science with pysr and symbolicregression.jl 667 (2023).
- 668 10. EO Voit, Perspective: Dimensions of the scientific method. *PLOS Comput. Biol*. **15**, e1007279 669 (2019).
- 670 11. TAA Adcock, PH Taylor, The physics of anomalous ('rogue') ocean waves. *Reports on Prog.* 671 *Phys*. **77**, 105901 (2014).
- 672 12. MA Tayfun, F Fedele, Wave-height distributions and nonlinear effects. *Ocean. Eng*. **34**, 673 1631–1649 (2007).
- 674 13. M Miche, Mouvements ondulatoires de la mer en profondeur constante ou decroissante. 675 *Annales de Ponts et Chaussees, 1944, pp(1) 26-78, (2)270-292, (3) 369-406* (1944).
- 676 14. J Gemmrich, C Garrett, Dynamical and statistical explanations of observed occurrence rates 677 of rogue waves. *Nat. Hazards Earth Syst. Sci*. **11**, 1437–1446 (2011).
- 678 15. F Fedele, J Brennan, S Ponce de León, J Dudley, F Dias, Real world ocean rogue waves 679 explained without the modulational instability. *Sci. Reports* **6**, 27715 (2016).
- 680 16. M Onorato, et al., Extreme waves, modulational instability and second order theory: wave 681 flume experiments on irregular waves. *Eur. J. Mech. - B/Fluids* **25**, 586–601 (2006).
- 682 17. K Trulsen, H Zeng, O Gramstad, Laboratory evidence of freak waves provoked by non-uniform
- 683 bathymetry. *Phys. Fluids* **24**, 097101 (2012) Publisher: American Institute of Physics. 684 18. JK Mallory, Abnormal Waves on the South East Coast of South Africa. *The Int. Hydrogr. Rev*. 685 (1974).
686 19. EG Dio
- For Exact To boosted occurrence rates
 DRAFTAI-01 Cool Control (Control Control Cont 19. EG Didenkulova, TG Talipova, EN Pelinovsky, Rogue Waves in the Drake Passage: Unpre-687 dictable Hazard in *Antarctic Peninsula Region of the Southern Ocean: Oceanography and* 688 *Ecology*, Advances in Polar Ecology, eds. EG Morozov, MV Flint, VA Spiridonov. (Springer International Publishing, Cham), pp. 101-114 (2021).
- 690 20. ML McAllister, S Draycott, TaA Adcock, PH Taylor, TSvd Bremer, Laboratory recreation of the 691 691 Draupner wave and the role of breaking in crossing seas. *J. Fluid Mech*. **860**, 767–786 (2019). 692 21. J Pearl, *Causality*. (Cambridge university press), (2009).
- 693 22. MA Tayfun, Distribution of large wave heights. *J. waterway, port, coastal, ocean engineering*
- 694 **116**, 686–707 (1990). 695 23. Y Goda, Reanalysis of Regular and Random Breaking Wave Statistics. *Coast. Eng. J*. **52**,
- 696 71–106 (2010).
697 24. T Tang. D Barra 24. T Tang, D Barratt, HB Bingham, TS van den Bremer, TA Adcock, The impact of removing the
- 698 high-frequency spectral tail on rogue wave statistics. *J. Fluid Mech*. **953**, A9 (2022).
- 699 25. DJ Korteweg, G De Vries, On the change of form of long waves advancing in a rectangular 700 canal, and on a new type of long stationary waves. *The London, Edinburgh, Dublin Philos.* 701 *Mag. J. Sci*. **39**, 422–443 (1895).
- 702 26. P Janssen, Shallow-water version of the Freak Wave Warning System, (ECMWF), Technical 703 memorandum 813 (2018).
- 704 27. PAEM Janssen, Nonlinear Four-Wave Interactions and Freak Waves. *J. Phys. Oceanogr*. **33**, 705 863–884 (2003) Publisher: American Meteorological Society.
- 706 28. F Ursell, The long-wave paradox in the theory of gravity waves. *Math. Proc. Camb. Philos.* 707 *Soc*. **49**, 685–694 (1953) Publisher: Cambridge University Press.
- 708 29. ECMWF, Part VII: ECMWF Wave model in *IFS Documentation CY47R3*, IFS Documentation. 709 (ECMWF), (2021).
- 710 30. K Hornik, Approximation capabilities of multilayer feedforward networks. *Neural Networks* **4**, 711 251–257 (1991).
- 712 31. V Nair, GE Hinton, Rectified linear units improve restricted boltzmann machines in *Proceedings* 713 *of the 27th International Conference on International Conference on Machine Learning*, 714 ICML'10. (Omnipress, Madison, WI, USA), pp. 807–814 (2010).
- 715 32. DP Kingma, J Ba, Adam: A method for stochastic optimization. *arXiv preprint arXiv:1412.6980* 716 (2014).
- 717 33. W Maddox, T Garipov, P Izmailov, D Vetrov, AG Wilson, A Simple Baseline for Bayesian 718 Uncertainty in Deep Learning. *arXiv:1902.02476 [cs, stat]* (2019) arXiv: 1902.02476.
- 719 34. J Peters, P Bühlmann, N Meinshausen, Causal inference by using invariant prediction: identifi-720 cation and confidence intervals. *J. Royal Stat. Soc. Ser. B (Statistical Methodol*. **78**, 947–1012 721 (2016).
- 722 35. J Peters, D Janzing, B Schölkopf, *Elements of Causal Inference: Foundations and Learning* 723 *Algorithms*, Adaptive Computation and Machine Learning series ed. F Bach. (MIT Press, Cambridge, MA, USA), (2017).
- 725 36. C Heinze-Deml, J Peters, N Meinshausen, Invariant causal prediction for nonlinear models. *J.* 726 *Causal Inference* **6**, 20170016 (2018).
- 727 37. P Xenopoulos, J Rulff, LG Nonato, B Barr, C Silva, Calibrate: Interactive analysis of probabilistic 728 model output. *IEEE Transactions on Vis. Comput. Graph*. **29**, 853–863 (2022).
- 729 38. MS Longuet-Higgins, On the statistical distribution of the height of sea waves. *JMR* **11**, 730 245–266 (1952).
- 731 39. M Schmidt, H Lipson, Distilling free-form natural laws from experimental data. *science* **324**,

81–85 (2009). 732

- 40. SM Udrescu, et al., Ai feynman 2.0: Pareto-optimal symbolic regression exploiting graph 733 modularity. *Adv. Neural Inf. Process. Syst*. **33**, 4860–4871 (2020). 734
- 41. SL Brunton, JL Proctor, JN Kutz, Discovering governing equations from data by sparse 735 identification of nonlinear dynamical systems. *Proc. national academy sciences* **113**, 3932– 736 3937 (2016). 737
- 42. KR Broløs, et al., An approach to symbolic regression using Feyn. *arXiv preprint* 738 *arXiv:2104.05417* (2021). 739
- 43. JH Holland, Genetic algorithms. *Sci. american* **267**, 66–73 (1992). 740
- 44. F Fedele, MA Tayfun, On nonlinear wave groups and crest statistics. *J. Fluid Mech*. **620**, 741 221–239 (2009) Publisher: Cambridge University Press. 742
- 45. F Fedele, J Herterich, A Tayfun, F Dias, Large nearshore storm waves off the Irish coast. *Sci.* 743 *Reports* **9**, 15406 (2019) Number: 1 Publisher: Nature Publishing Group. 744
- 46. O Gramstad, K Trulsen, Influence of crest and group length on the occurrence of freak waves. 745 *J. Fluid Mech*. **582**, 463–472 (2007) Publisher: Cambridge University Press. 746
- 47. W Xiao, Y Liu, G Wu, DKP Yue, Rogue wave occurrence and dynamics by direct simulations 747 of nonlinear wave-field evolution. *J. Fluid Mech*. **720**, 357–392 (2013) Publisher: Cambridge 748 University Press. 749
- 48. J Gemmrich, J Thomson, Observations of the shape and group dynamics of rogue waves. 750 *Geophys. Res. Lett*. **44**, 1823–1830 (2017). 751
- 49. DW Apley, J Zhu, Visualizing the Effects of Predictor Variables in Black Box Supervised 752 Learning Models. *arXiv:1612.08468 [stat]* (2019) arXiv: 1612.08468. 753
- 50. N Mori, M Onorato, PA Janssen, On the estimation of the kurtosis in directional sea states for 754 freak wave forecasting. *J. Phys. Oceanogr*. **41**, 1484–1497 (2011). 755
- 51. LH Ying, Z Zhuang, EJ Heller, L Kaplan, Linear and nonlinear rogue wave statistics in the 756 presence of random currents. *Nonlinearity* **24**, R67–R87 (2011) Publisher: IOP Publishing. 757
- 52. M Casas-Prat, LH Holthuijsen, Short-term statistics of waves observed in deep water. *J.* 758 *Geophys. Res. Ocean*. **115** (2010). 759
- 53. J Gemmrich, L Cicon, Generation mechanism and prediction of an observed extreme rogue 760 wave. *Sci. Reports* **12**, 1–10 (2022). 761
- 54. IR Young, The determination of confidence limits associated with estimates of the spectral 762 peak frequency. *Ocean. Eng*. **22**, 669–686 (1995). 763
- 55. JD Fenton, The numerical solution of steady water wave problems. *Comput. & Geosci*. **14**, 764 357–368 (1988). 765
- 56. M Serio, M Onorato, A R a Osborne, P Janssen, On the computation of the Benjamin-Feir 766 Index. *Nuovo Cimento della Soc. Italiana di Fisica C* **28**, 893–903 (2005). 767
- 57. J Bradbury, et al., JAX: composable transformations of Python+NumPy programs (2018). 768
- 58. J Heek, et al., Flax: A neural network library and ecosystem for JAX (2020).
59. M Hessel, et al., Optax: composable gradient transformation and optimisation, in JAX! (2020). 770
- M Hessel, et al., Optax: composable gradient transformation and optimisation, in JAX! (2020).
- 60. N Mori, PAEM Janssen, On Kurtosis and Occurrence Probability of Freak Waves. *J. Phys.* 771 *Oceanogr*. **36**, 1471–1483 (2006) Publisher: American Meteorological Society. 772
- 61. D Häfner, Ph.D. thesis (2022). 773
- 62. F Pedregosa, et al., Scikit-learn: Machine learning in Python. *J. Mach. Learn. Res*. **12**, 774 2825–2830 (2011). 775
- 63. D Jomar, PyALE: A Python implementation of accumulated local effect plots (2020). 776
- 64. CR Harris, et al., Array programming with NumPy. *Nature* **585**, 357–362 (2020). 777
- 65. P Virtanen, et al., SciPy 1.0: Fundamental Algorithms for Scientific Computing in Python. *Nat.* 778 *Methods* **17**, 261–272 (2020). 779
- 66. JD Hunter, Matplotlib: A 2d graphics environment. *Comput. Sci. & Eng*. **9**, 90–95 (2007). 780
- 67. ML Waskom, seaborn: statistical data visualization. *J. Open Source Softw*. **6**, 3021 (2021). 781 68. Wes McKinney, Data Structures for Statistical Computing in Python in *Proceedings of the 9th* 782
- *Python in Science Conference*, eds. Stéfan van der Walt, Jarrod Millman. pp. 56 61 (2010). 783
- 69. T Kluyver, et al., Jupyter notebooks - a publishing format for reproducible computational 784 workflows in *Positioning and Power in Academic Publishing: Players, Agents and Agendas*, 785 eds. F Loizides, B Scmidt. (IOS Press, Netherlands), pp. 87–90 (2016). The match of the matc